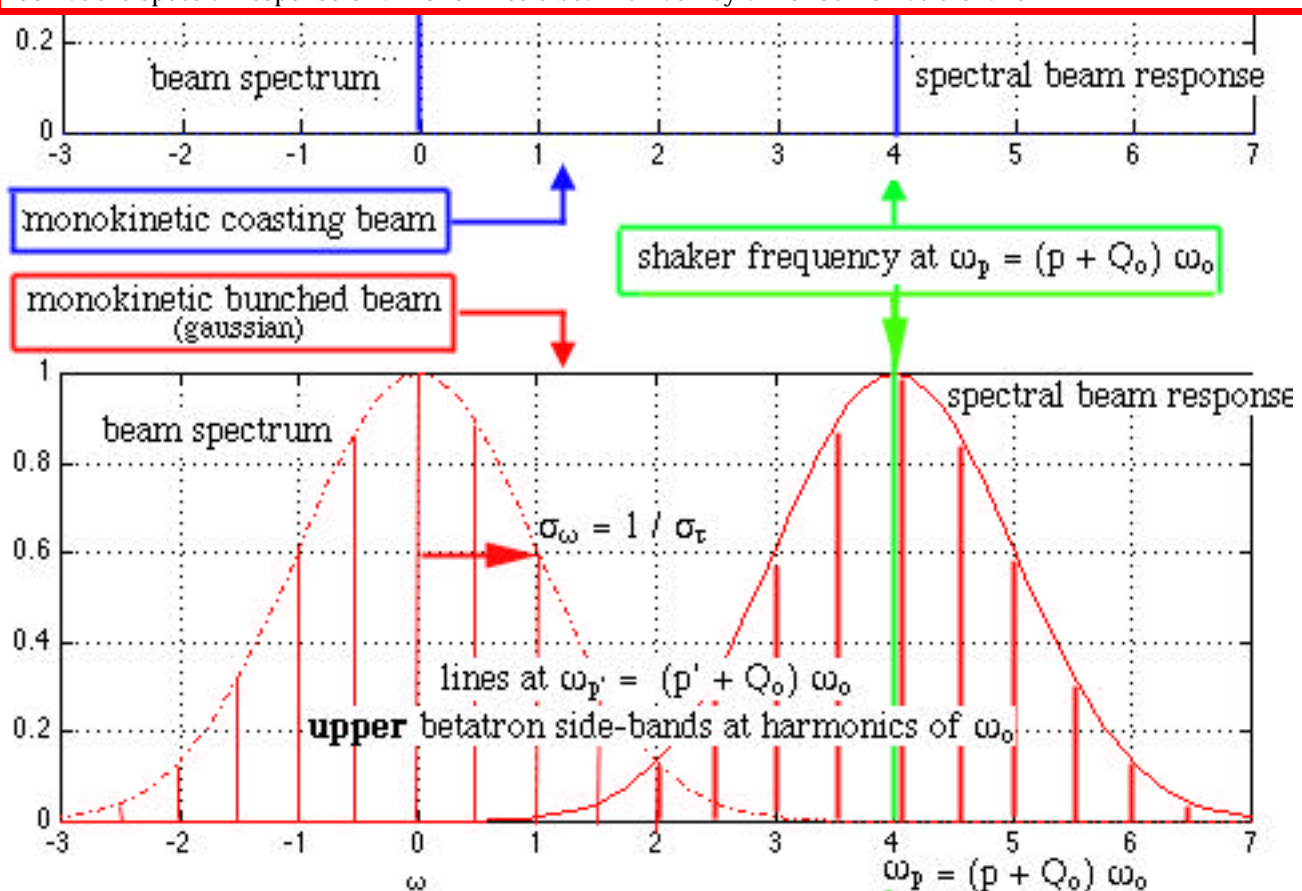


When the ESRF works with a high chromaticity, the intensity in single bunch mode is suspected to be limited by a fast transverse instability of "post head-tail" type (with growing time shorter than a synchrotron period) . Elements of a theory* in frequency space are presented , using J.L. Laclare 's formalism.

The work on this topic is part of a thesis by Ph.Kernel: some aspects are recent, not fully achieved and possibly need more reflection .

Look at the spectral response of a monokinetic beam driven by a monochromatic shaker



The spectral response is the beam spectrum shifted to the shaker frequency at ω_p

J.L. Laclare's formalism is more easily understood when we define the matrix elements:

$\mathbf{A}_{p', p}$ = the amplitude frequency line $\sigma(\omega_{p'})$ in response to a unit excitation at ω_p
 = the spectrum of the longitudinal bunch density shifted at ω_p
 (later designed as a "shaker mode" spectrum , centered at ω_p)
 $= \delta_{p, p'}$ for a **coasting beam** , $= \exp - [(\omega_p - \omega_{p'})^2 \sigma_t^2 / 2]$ for a **gaussian beam** .

such a matrix **A** works like the "impulse response" of a filter in frequency space

* See also a

previous and different analysis by R.D.Ruth and J.M. Wang: Vertical fast blow-up in a single bunch. IEEE Transactions on Nuclear Sciences, NS-28, N°3, June 1982.

An approximative criterion for the stability of an intense bunched beam at high chromaticity .

part 1(monokinetic)

page 2.

Transverse modes of a monokinetic bunch , broad band impedance

the stability of a collective bunch oscillation is given by
the imaginary part of the complex frequency shift $\Delta\omega_c^0$ from $Q_o \omega_o$

narrow band impedance $e^* Z_T(\omega)$

1)admit the **coasting beam** result :

$$\Delta\omega_{cp}^0 = \Lambda j Z_T(\omega_p) , \quad \Lambda = \frac{c I_o}{4\pi Q_o E_o / e}$$

[or $\Delta\omega_{cp}^0 \sigma(\omega_p) = \Lambda \delta_{p',p} j Z_T(\omega_p) \sigma(\omega_p)$]

2)deduce the **bunched beam** formula :

$$\Delta\omega_c^0 \sigma(\omega_p) = \Lambda \mathbf{A}_{p',p} j Z_T(\omega_p) \sigma(\omega_p)$$

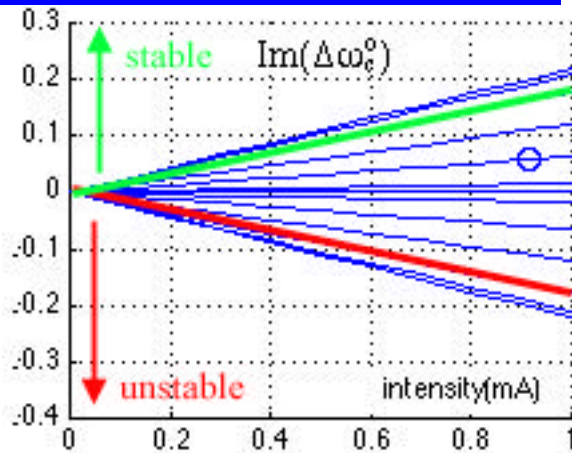
(only one input line)

broad band impedance $Z_T(\omega)$

3)same **coasting beam** formula :

$$\Delta\omega_{cp}^0 = \Lambda j Z_T(\omega_p) , \quad \Lambda = \frac{c I_o}{4\pi Q_o E_o / e}$$

stability : upper side band of
mode " p " : $\omega_{cp}^+ = (p + Q_o + \Delta\omega_{cp}^0) \omega_o$
unstable if $\text{Re}(Z_T(\omega_{cp}^+)) < 0$.

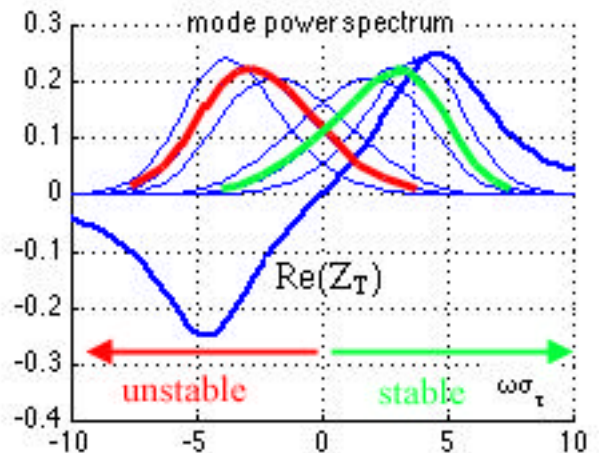


4)**bunched beam** : sum all input lines

$$\Delta\omega_c^0 \sigma(\omega_p) = \Lambda \bullet_p j Z_T(\omega_p) A_{p',p} \sigma(\omega_p)$$

stability :

1- look for **computed solutions**
of this standard eigen-value problem :



stability : 2 - look for **approximative solutions** of the eigen-value problem :
 $\sigma_q(\omega_p)$ an eigen-mode spectrum associated with an eigen-value $\Delta\omega_{cq}^0$,

C_q the associated eigen-value of matrix \mathbf{A} , the n^{**} :

$$\Delta\omega_{cq}^0 = \Lambda j [Z_{Tq}]_{\text{eff}} C_q, \text{ with effective impedance : } [Z_{Tq}]_{\text{eff}} = \frac{\bullet_p Z_T(\omega_p) \sigma_q^2(\omega_p)}{\bullet_p \sigma_q^2(\omega_p)}$$

Approximation of σ_q by a "shaker mode" : $\sigma_q(\omega_p) = \exp -[(\omega_p - \omega_q)^2 \sigma_\tau^2 / 2]$, then $C_q = \frac{2\sqrt{\pi/3}}{\omega_o \sigma_\tau}$ (indpt of q) .

Conclusion : using J.L.Laclare 's formalism, we recover rapidly the following property :

Adding the contributions of upper and lower side-bands : monokinetic coasting and bunched beams are unstable with respect to any impedance which has a resistive component. The frequency shift $\text{Re}(\Delta\omega_{cq})$ and growing time $\text{Im}(\Delta\omega_{cq})^{-1}$ of any mode "q" can be computed

* by analogy between a shaker tuned at ω_p and a narrow band impedance $Z_T(\omega)$
which overlaps only one line ω_p of the transverse returned signal.

** F.J. Sacherer derived such approximate but very useful formulas for the stability of head-tail sine modes